

Universitat Politècnica de Catalunya, Barcelona, 15 – 19 April 2013

Numerical Simulation of Fast Transient Phenomena in Fluid-Structure Systems

A Short Course by F. Casadei

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The Joint Research Centre (JRC)

7 Institutes in 5 Member States:



IRMM – Geel, Belgium

- Institute for Reference Materials and Measurements



IE – Petten, The Netherlands

- Institute for Energy



ITU – Karlsruhe, Germany

- Institute for Transuranium elements



IPSC - IHCP - IES – Ispra, Italy

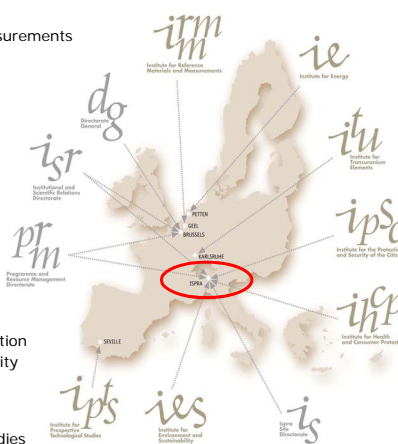
- **Institute for the Protection and
Security of the Citizen**

- Institute for Health and Consumer Protection
- Institute for Environment and Sustainability



IPTS – Seville, Spain

- Institute for Prospective Technological Studies



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The European Laboratory for Structural Assessment (**ELSA**)



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ELSA **Facilities:**

Large Hopkinson Bar
("Fast" Dynamics)



Reaction
Wall
("Slow"
Dynamics)



“Slow” dynamics: earthquake ...

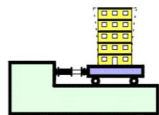


Kobe (Japan) earthquake. January 17, 1995 (6,300 fatalities)

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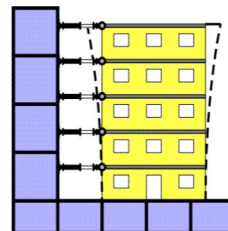
Earthquake testing of structures: two complementary methods

Shaking Table



- Dynamic test (real time)
- Small-scale models

Reaction Wall



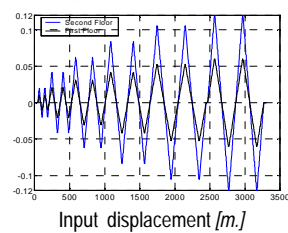
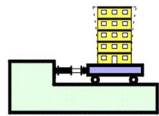
- Pseudo-dynamic test (expanded time scale)
- Full-scale models

(ELSA)

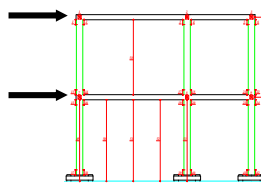
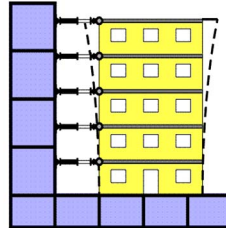
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Imposed motion

Shaking Table Base motion



Reaction Wall Floors motion

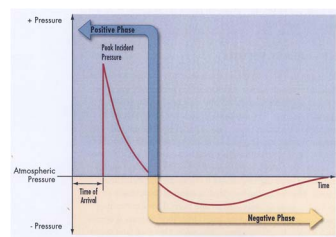


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"Fast" dynamics: impact, explosion ...



Twin towers. September 11, 2001
(2,967 fatalities)

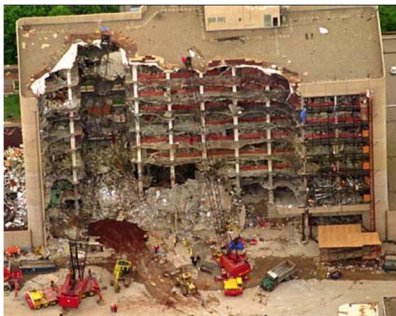


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Progressive Collapse: a "slow" phenomenon often initiated by a "fast" transient



Twin towers.
Structure collapse



Murrah Building, Oklahoma City.
April 19, 1995 (168 fatalities)

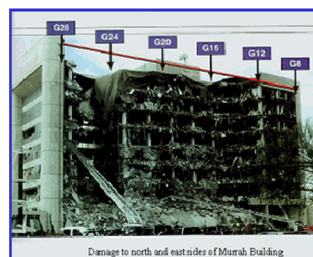


Aladdin Hotel, Las Vegas. April
27, 1998 (controlled demolition)

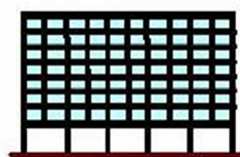
Example of Progressive Collapse due to terrorist attack



Murrah Building prior to blast



Damage to north and east sides of Murrah Building



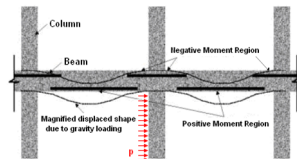
Portion progressively collapsed

Column G20
blasted away

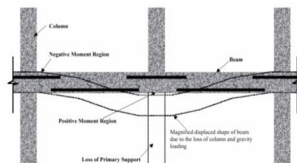
Shear failure of
adjacent columns

Blast Simulator (Under construction)

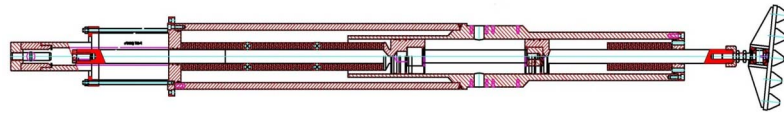
Design for Gravity Loads



Response to Column Loss



(Courtesy of San Diego University)



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Contents

- I – Introduction (Structures) ←
- II – ALE formulation (Fluids)
- III – Classical Fluid-Structure Interaction
- IV – Advanced FSI (Failure/Fragmentation)
- V – Further topics and applications

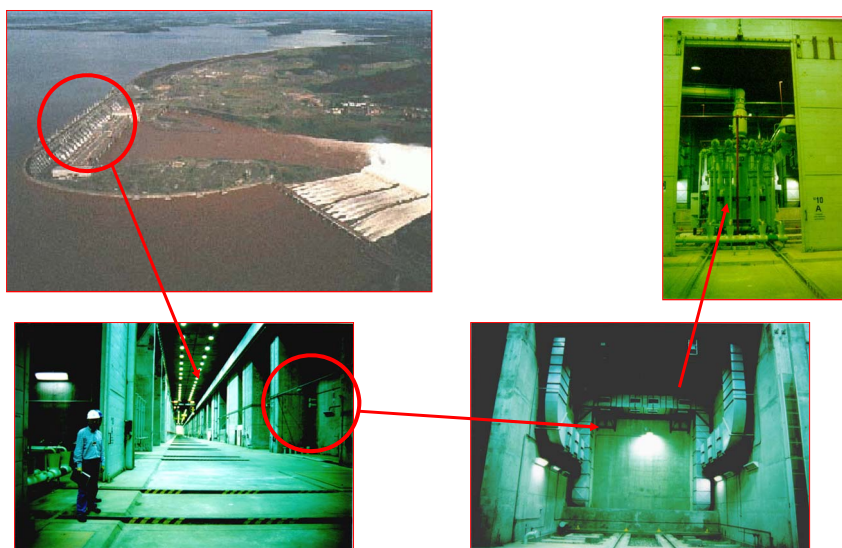
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Credits & Acknowledgments

- Structural Dynamics team at JRC Ispra (70's – to date):
 - J. Donea, J.P. Halleux, S. Giuliani, F. Casadei, M. Larcher, G. Solomos, G. Giannopoulos, G. Valsamos, N. Leconte ...
- Structural Dynamics team at CEA Saclay (70's – to date):
 - H. Bung, P. Galon, M. Lepareux, V. Faucher, A. Beccantini ...
- Contributions from research/academic bodies and industry:
 - UPC (A. Huerta, P. Diez, F. Verdugo), Cachan, CRS4, Lyon, ...
 - EDF, Onera, ENEL, Snecma, SamTech ...
- Three decades of code development:
 - EURDYN, Castem-PLEXUS / PLEXIS-3C, EUROPLEXUS ...
 - Commercial / Academic versions now available from CEA

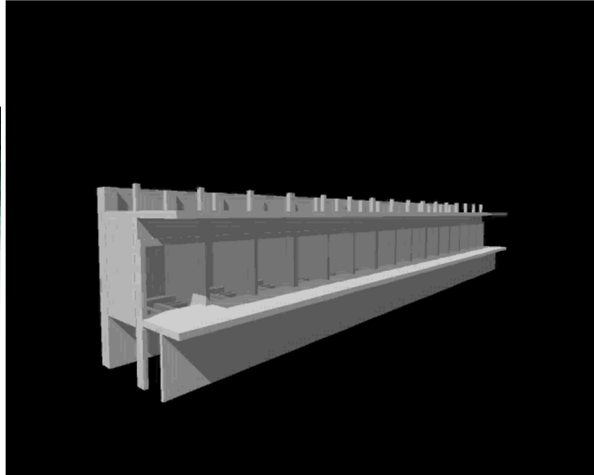
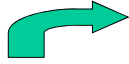
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Introductory Example (Courtesy of ENEL-Hydro)



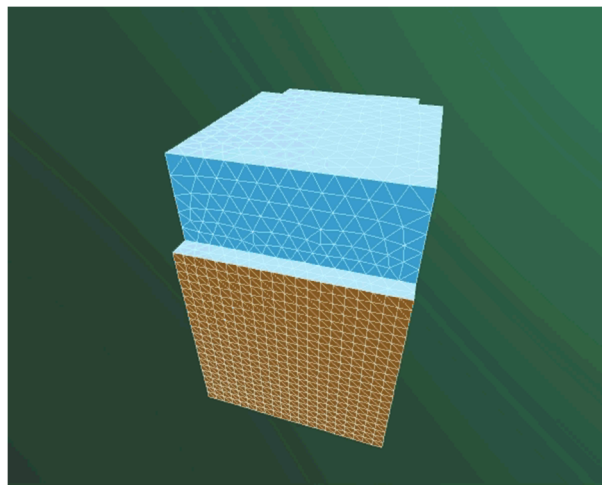
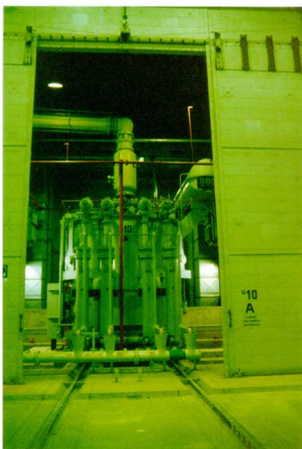
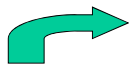
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Introductory Example (2)



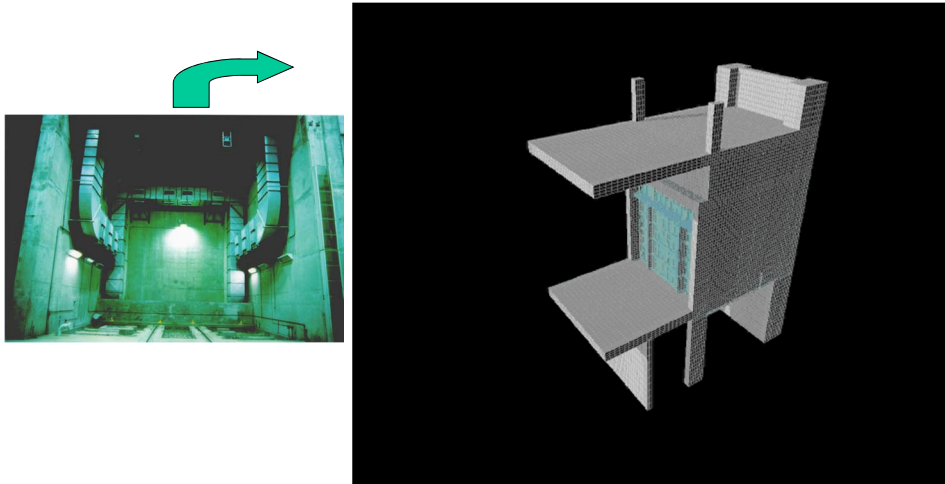
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Introductory Example (3)



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Introductory Example (4)



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Application Spectrum

- **Energy (safety issues):** nuclear and fossil-fueled plants, electrical devices, chemical plants, pressure vessels, ...
- **Civil engineering:** earthquakes, soil-structure interactions, building vulnerability to terrorist attacks ...
- **Marine/Offshore:** ships and submarines, oil industry, pipelines, cables ...
- **Transportation:** crash, road barriers, tunnels safety, trains/metro stations and rolling stock...
- etc...

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Goals/Characteristics

- **Simulate fast transient dynamic phenomena:** explosions, crashes, impacts ...
- **Short time scale:** (typically *ms* to a few *s*) with large frequencies spectrum
- **Geometric and material non-linearities:** large motions, large strains, plasticity, visco-plasticity, damage ...
- **Structures and fluids:** heterogeneity, interaction phenomena, ...
- **For reliable solutions,** simple and robust numerical methods are needed: direct time integration, explicit schemes ...

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Detailed Contents of Part I

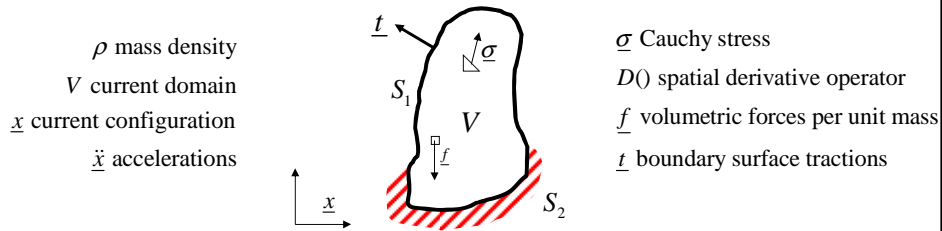
- **Introductory example of FSI problem**
 - ✓ Application spectrum and goals
- **Modeling the structural domain**
 - ✓ Equilibrium equations
 - ✓ Explicit time integration scheme
- **Treatment of essential boundary conditions**
 - ✓ The Lagrange multipliers method
 - ✓ Solving the linear system

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Computational Framework

- Governing equation for structural domain: principle of virtual work (conservation of momentum, i.e. equilibrium in a dynamic sense)

$$\int_V \rho \ddot{\underline{x}} \delta \underline{x} dV + \int_V \underline{\sigma} D(\delta \underline{x}) dV - \int_V \rho \underline{f} \delta \underline{x} dV - \int_{S_1} \underline{t} \delta \underline{x} dS = 0$$

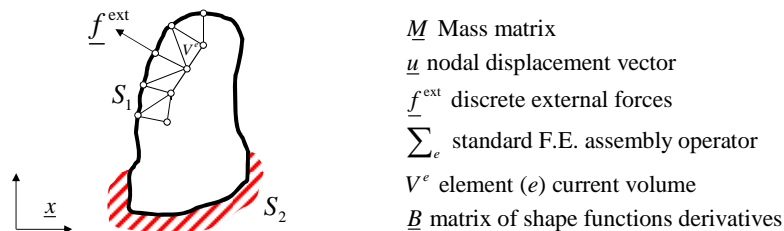


Must hold for all variations $\delta \underline{x}$ of configuration (virtual displacements) compatible with essential b.c.s on S_2 . 21

Computational Framework (2)

- This integral form lends itself to direct application of F.E. method. Upon spatial discretization:

$$\underline{M} \ddot{\underline{u}} = \underline{f}^{\text{ext}} - \sum_e \int_{V^e} \underline{B}^T \underline{\sigma} dV$$



This set of discrete differential equations in time is decoupled by diagonalization (lumping) of mass matrix \underline{M} 22

$$\underline{M}\ddot{\underline{u}} = \underline{f}^{\text{ext}} - \sum_e \int_{V^e} \underline{B}^T \underline{\sigma} dV$$

Computational Framework (3)

- **Description is Lagrangian:** nodes and G.P.s remain always associated with same material point (particle)
- **Stress is “true”:** expressed in fixed reference (but corotational formulation can be useful for beams/plates/shells)
- **All RHS terms are known ($\underline{f}^{\text{ext}}, \underline{B}$) or computable:** stresses must be obtained via material constitutive law
- **Diagonalization of \underline{M} by lumping:** $\underline{M}^e = \int_{V^e} \underline{N} \rho dV$
where \underline{N} are the element shape functions
- **We work on current configuration:** no need to define a reference configuration and no use of (total) deformation

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$$\underline{M}\ddot{\underline{u}} = \underline{f}^{\text{ext}} - \sum_e \int_{V^e} \underline{B}^T \underline{\sigma} dV$$

Direct Time Integration

- Time integration is achieved via the Central Difference scheme, usually written as:

$$\begin{aligned} \underline{\dot{u}}^{n+1} &= \underline{\dot{u}}^n + \frac{\Delta t}{2} (\ddot{u}^n + \ddot{u}^{n+1}) \\ \underline{u}^{n+1} &= \underline{u}^n + \Delta t (\underline{\dot{u}}^n + \frac{\Delta t}{2} \ddot{u}^n) \end{aligned}$$

n stays for time t^n

$n+1$ stays for $t^{n+1} = t^n + \Delta t$

Δt is the time increment

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$$\begin{aligned}\dot{\underline{u}}^{n+1} &= \dot{\underline{u}}^n + \frac{\Delta t}{2}(\ddot{\underline{u}}^n + \ddot{\underline{u}}^{n+1}) \\ \underline{u}^{n+1} &= \underline{u}^n + \Delta t(\dot{\underline{u}}^n + \frac{\Delta t}{2}\ddot{\underline{u}}^n)\end{aligned}$$

Direct Time Integration (2)

- These formulas are a particularization of the well-known Newmark integration formulas:

$$\dot{\underline{u}}^{n+1} = \dot{\underline{u}}^n + \Delta t[(1-\gamma)\ddot{\underline{u}}^n + \gamma\ddot{\underline{u}}^{n+1}]$$

$$\underline{u}^{n+1} = \underline{u}^n + \Delta t\dot{\underline{u}}^n + \frac{\Delta t^2}{2}[(1-2\beta)\ddot{\underline{u}}^n + 2\beta\ddot{\underline{u}}^{n+1}]$$

written for $\gamma = 1/2$ and $\beta = 0$.

$$M\ddot{\underline{u}} = \underline{f}^{\text{ext}} - \underline{f}^{\text{int}}$$

- These two equations, plus the equilibrium, can be solved for \underline{u} , $\dot{\underline{u}}$, $\ddot{\underline{u}}$ upon step-by-step marching in time.
- This particular choice for β renders the scheme explicit, while the chosen γ ensures no numerical damping.

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$$\begin{aligned}\dot{\underline{u}}^{n+1} &= \dot{\underline{u}}^n + \frac{\Delta t}{2}(\ddot{\underline{u}}^n + \ddot{\underline{u}}^{n+1}) \\ \underline{u}^{n+1} &= \underline{u}^n + \Delta t(\dot{\underline{u}}^n + \frac{\Delta t}{2}\ddot{\underline{u}}^n)\end{aligned}$$

Direct Time Integration (3)

How is the scheme used in practice?

- Introduce a mid-step velocity:

$$\underline{v}^{n+1/2} \doteq \dot{\underline{u}}^n + \frac{\Delta t}{2}\ddot{\underline{u}}^n$$

which transforms configuration n into $n+1$ over the step.

- The second equation becomes:

$$\underline{u}^{n+1} = \underline{u}^n + \Delta t \cdot \underline{v}^{n+1/2}$$

- Carry on mid-step velocities rather than full-step ones. The first equation becomes:

$$\underline{v}^{n+3/2} = \underline{v}^{n+1/2} + \Delta t \cdot \ddot{\underline{u}}^{n+1}$$

$$\underline{v}^{n+3/2} \doteq \dot{\underline{u}}^{n+1} + \frac{\Delta t}{2}\ddot{\underline{u}}^{n+1} = \dot{\underline{u}}^n + \frac{\Delta t}{2}\ddot{\underline{u}}^n + \frac{\Delta t}{2}\ddot{\underline{u}}^{n+1} + \frac{\Delta t}{2}\ddot{\underline{u}}^{n+1} = \underline{v}^{n+1/2} + \Delta t \cdot \ddot{\underline{u}}^{n+1}$$

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Direct Time Integration (4)

- The final algorithm reads:

$$\underline{u}^{n+1} = \underline{u}^n + \Delta t \cdot \underline{v}^{n+1/2}$$

$$\ddot{\underline{u}}^{n+1} = \underline{M}^{-1}(\underline{f}^{\text{ext}(n+1)} - \sum_e \int_{V^e(n+1)} \underline{B}^T \underline{\sigma}^{e(n+1)} dV)$$

$$\underline{v}^{n+3/2} = \underline{v}^{n+1/2} + \Delta t \cdot \ddot{\underline{u}}^{n+1}$$

A new configuration is obtained first. On this known configuration, equilibrium is enforced. The new mid-step velocity is obtained last.

- This scheme is explicit.
- If Δt varies in time, the only change is in the third equation, which becomes:

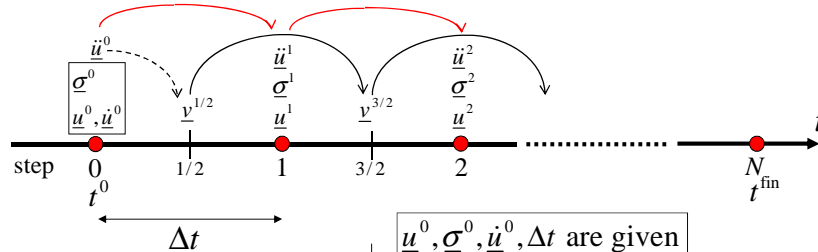
$$\underline{v}^{n+3/2} = \underline{v}^{n+1/2} + \frac{\Delta t^n + \Delta t^{n+1}}{2} \cdot \ddot{\underline{u}}^{n+1}$$

with:

$$\begin{aligned} \Delta t^n &\equiv t^{n+1} - t^n \\ \Delta t^{n+1} &\equiv t^{n+2} - t^{n+1} \end{aligned} \quad 27$$

How does one obtain $\underline{\sigma}^{e(n+1)}$? (see below)

Scheme start-up and marching



$$\begin{aligned} \underline{u}^{n+1} &= \underline{u}^n + \Delta t \cdot \underline{v}^{n+1/2} \\ \ddot{\underline{u}}^{n+1} &= \underline{M}^{-1}(\underline{f}^{\text{ext}} - \underline{f}^{\text{int}})^{n+1} \\ \underline{v}^{n+3/2} &= \underline{v}^{n+1/2} + \Delta t \cdot \ddot{\underline{u}}^{n+1} \end{aligned}$$

$\underline{u}^0, \underline{\sigma}^0, \underline{\dot{u}}^0, \Delta t$ are given

$$n = -1$$

$$\ddot{\underline{u}}^0 = \underline{M}^{-1}(\underline{f}^{\text{ext}} - \underline{f}^{\text{int}})^0$$

$$\underline{v}^{1/2} = \underline{\dot{u}}^0 + (\Delta t / 2) \cdot \ddot{\underline{u}}^0$$

$$n \leftarrow n + 1$$

$$\underline{u}^1 = \underline{u}^0 + \Delta t \cdot \underline{v}^{1/2}$$

$$\ddot{\underline{u}}^1 = \underline{M}^{-1}(\underline{f}^{\text{ext}} - \underline{f}^{\text{int}})^1$$

$$\underline{v}^{3/2} = \underline{v}^{1/2} + \Delta t \cdot \ddot{\underline{u}}^1$$

etc.

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Scheme start-up and marching (2)

- For practical reasons, the code also computes the full-step velocities:

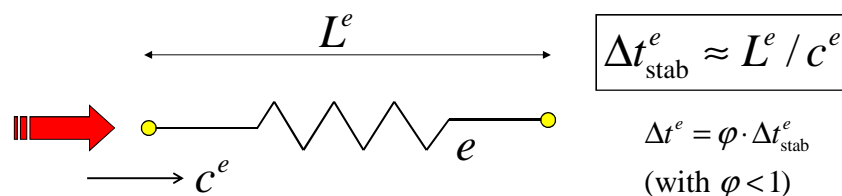
$$\underline{\dot{u}}^{n+1} = \underline{v}^{n+1/2} + \frac{\Delta t}{2} \underline{\ddot{u}}^{n+1}$$

- These are the velocities printed out in the listing and visualized in post-processing
- However, the fundamental quantity in the time integration scheme is the mid-step velocity!

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Integration Scheme Characteristics

- Central difference scheme is second-order accurate and introduces no numerical damping
- However, it is conditionally stable (Courant):



- In highly non-linear cases, small steps are needed even with unconditionally stable schemes, to get good accuracy

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Integration Scheme Characteristics (2)

- Spectral analysis shows that the central difference scheme tends to produce frequencies slightly higher than physical ones. Same effect is obtained using a consistent mass matrix.
- However, use of a lumped mass matrix tends to reduce frequency values.
- Therefore, combination of CD time integrator with a lumped mass matrix gives optimal numerical precision.
- This is a remarkable result, since the final equations are completely decoupled: contrary to classical FE method, there are no matrices to assemble and no need for system solvers (except for treatment of essential BCs).

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Integration Scheme Characteristics (3)

- See: S.W.Key, “*Transient Response by Time Integration*”.

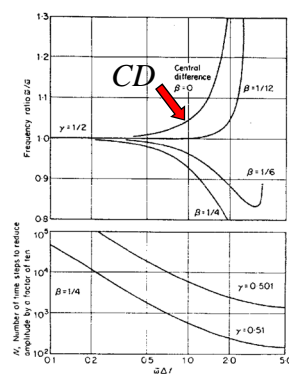


FIG. 6. The frequency response for $\gamma = 0.5$ and damping for $\beta = 0.25$ for the Newmark β time integrator.

Scheme effect on frequency

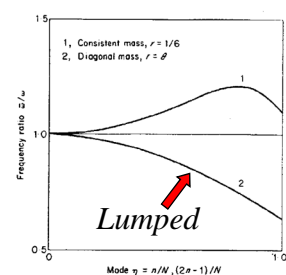


FIG. 2. The ratio of approximate frequency, $\tilde{\omega}$, to true frequency, ω , versus non-dimensional wave number for a piecewise linear spatial approximation.

Mass matrix effect on frequency

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Stress Update

- To solve the equilibrium equation for the new accelerations:

$$\ddot{\underline{u}}^{n+1} = \underline{M}^{-1}(\underline{f}^{\text{ext}(n+1)} - \sum_e \int_{V^{e(n+1)}} \underline{B}^T \underline{\sigma}^{e(n+1)} dV)$$

one needs the new stress $\underline{\sigma}^{e(n+1)}$.

- In general one can formally write:

$$\underline{\sigma}^{n+1} = \underline{\sigma}^n + \Delta \underline{\sigma}$$

$$\Delta \underline{\sigma} = H(\underline{\sigma}^n, \Delta \underline{\epsilon}, \underline{p}, \dot{\underline{\epsilon}}, \dots) \quad (\text{Rate form})$$

H Constitutive law

$\Delta \underline{\sigma}$ stress increment over the step \underline{p} hardening parameters (e.g. plasticity)

$\Delta \underline{\epsilon}$ strain increment over the step $\dot{\underline{\epsilon}}$ strain rate (e.g. viscous behaviour)

- Note that the total deformation $\underline{\epsilon}$ does not appear anywhere and is not used in the process.

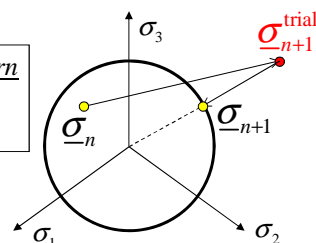
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Elasto-plastic material

As an example of non-linear material behaviour consider the important case of metal plasticity:

- Rate-independent deviatoric plasticity model with Von Mises yield criterion:

*Radial return
method
(Wilkins)*



- “Trial” stress (elastic):

$$\underline{\sigma}_{n+1}^{\text{trial}} = \underline{\sigma}_n + \underline{C} \cdot \Delta \underline{\epsilon}$$

- If trial stress lies outside yield surface, perform radial return onto (current) yield surface. No iterations!

How does one compute $\Delta \underline{\epsilon}$ from displacement increments (or velocities) in the presence of *geometrical non-linearities* (large strains and large motions, in particular *large rotations*?)

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Geometric non-linearities

A large-displacement large-strain formulation is adopted for full generality. For continuum-like FE:

- Compute spatial velocity gradient:

$$\underline{L} = \partial \dot{\underline{x}} / \partial \underline{x}$$

- Use additive decomposition to separate instantaneous deformation (symmetric) from rotation (antisymmetric part):

$$\underline{L} = \underline{D} + \underline{W}$$

$$\underline{D} = \frac{1}{2}(\underline{L} + \underline{L}^T) \text{ (stretching i.e. rate of deformation)}$$

$$\underline{W} = \frac{1}{2}(\underline{L} - \underline{L}^T) \text{ (spin i.e. rate of rotation)}$$

- We obtain then:

$$\dot{\underline{\epsilon}} = \underline{D} \quad ; \quad \Delta \underline{\epsilon} = \underline{D} \cdot \Delta t$$

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Geometric non-linearities (2)

For a continuum element the state of stress of interest to us, Cauchy stress $\underline{\sigma}$, is referred to a fixed frame in space. Consequently, its time derivative is not invariant with respect to rotation: $\dot{\underline{\sigma}}$ is not objective.

- An objective rate of stress $\hat{\underline{\sigma}}$ can be obtained under the form: $\hat{\underline{\sigma}} = \dot{\underline{\sigma}} - \underline{A}\underline{\sigma} + \underline{\sigma}\underline{A}$ where \underline{A} is an appropriate vorticity matrix.

- In the Zaremba-Jaumann-Noll formulation: $\underline{A} \doteq \underline{W}$ (other choices are possible, e.g. Green-Naghdi).

- However, the above considerations are valid only in an infinitesimal sense, while we need to use finite increments.

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Geometric non-linearities (3)

Set up following incrementally objective scheme to update the Cauchy stress (2D case for simplicity) in three phases:

- Let α be the angle of rotation over Δt and let $\theta = \alpha / 2$

$$\tan \theta = (\Delta t / 2) \cdot W_{12}^{n+1/2}$$

1. Apply first half of the rotation increment:

$$\underline{\sigma}^{n*} = \underline{R} \underline{\sigma}^n \underline{R}^T \quad \text{with} \quad \underline{R} = \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix}$$

2. Apply the constitutive equation:

$$\underline{\sigma}^{(n+1)*} = \underline{\sigma}^{n*} + \underline{\underline{C}} \cdot \Delta t \cdot \underline{D}^{n+1/2}$$

3. Apply second half of the rotation increment:

$$\underline{\sigma}^{n+1} = \underline{R} \underline{\sigma}^{(n+1)*} \underline{R}^T$$

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Geometric non-linearities (4)

For structural elements (bars, beams, shells) use co-rotational formulation:

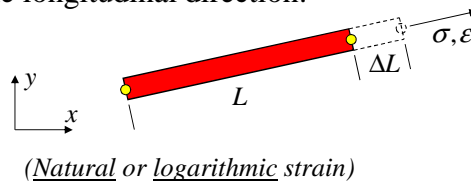
- The stress is measured in a reference frame that rotates with the element
- This greatly simplifies the stress increment procedure: the stress can be incremented directly by applying the constitutive law.

Example (bar element). In the longitudinal direction:

$$\Delta \varepsilon = \frac{\Delta L}{L}$$

$$\rightarrow \Delta \sigma$$

$$\sum \Delta \varepsilon = \sum \frac{\Delta L}{L} = \int_{L_0}^L \frac{dL}{L} = \ln \frac{L}{L_0}$$



A small-strain formulation would be:

$$\Delta \varepsilon = \frac{\Delta L}{L_0}$$

$$\rightarrow$$

$$\sum \Delta \varepsilon = \frac{L - L_0}{L_0}$$

(Engineering strain)

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Advantages of the Method

- Transient dynamic problem: find $\underline{\sigma}$ on new configuration \underline{x} (known) from $\underline{\sigma}^{\text{old}}$ and deformation process between $\underline{x}^{\text{old}}$ and \underline{x}
- Compare implicit methods: find $\underline{\sigma}$ and \underline{x} simultaneously, typically by iterative procedures and convergence criteria.
- The proposed method is particularly simple for complex non-linear problems, hence very robust.
- Direct application of virtual work principle plus second-order accurate time integration scheme, guarantee high accuracy of numerical results.

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Checking the Solution Quality

- The quality of the obtained solution can be checked globally by computing at each time step the energy balance.
- Initially, set: $W_0^{\text{ext}} \doteq E_0^{\text{int}} + E_0^{\text{kin}}$
- At any time, the balance error can be computed as:

$$\mathcal{E} = \frac{W^{\text{ext}} - (E^{\text{int}} + E^{\text{kin}})}{W^{\text{ext}}} \quad \text{or perhaps better:} \quad \mathcal{E} = \frac{W^{\text{ext}} - (E^{\text{int}} + E^{\text{kin}})}{\max(|W^{\text{ext}}|, |W_0^{\text{ext}}|)}$$

- This error indicator is used *a posteriori* in order to check the previously obtained solution and must not be confused with convergence parameters typical of iterative approaches

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Essential Boundary Conditions

Essential conditions are imposed via Lagrange multipliers.

- Assume a linear set of constraints on the velocities:

$$\underline{C}\underline{v} = \underline{b}$$

- Both \underline{C} and \underline{b} are known, and can be function of time.
- The equilibrium equations for the subset of d.o.f.s concerned become, introducing unknown reactions \underline{r} :

$$\underline{m}\underline{a} = \underline{f}^e - \underline{f}^i + \underline{r}$$

- Without loss of generality, the unknown reactions can be expressed via a vector $\underline{\lambda}$ of Lagrange multipliers:

$$\underline{r} = \underline{C}^T \underline{\lambda}$$

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Finding the Lagrange Multipliers

- Replacing into the equilibrium equations yields:

$$\underline{m}\underline{a} = \underline{f}^e - \underline{f}^i + \underline{C}^T \underline{\lambda}$$

- Multiplying both members by $\underline{C}\underline{m}^{-1}$ gives:

$$\underline{C}\underline{a} = \underline{C}\underline{m}^{-1}(\underline{f}^e - \underline{f}^i) + \underbrace{\underline{C}\underline{m}^{-1}\underline{C}^T}_{\underline{B}^*} \underline{\lambda}$$

Matrix of connections

- The Lagrange multipliers are obtained symbolically from:

$$\underline{B}^* \underline{\lambda} = \underline{C}\underline{a} - \underline{C}\underline{m}^{-1}(\underline{f}^e - \underline{f}^i)$$

- To obtain $\underline{\lambda}$, we must first express the term $\underline{C}\underline{a}$ as a function of known quantities, by using the constraint and the time integration scheme.

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Finding the Lagrange Multipliers (2)

- The CD scheme for the velocity and constant Δt is:

$$\underline{v}^{n+3/2} = \underline{v}^{n+1/2} + \Delta t \cdot \underline{a}^{n+1}$$

- Substituting this into the constraint $\underline{Cv} = \underline{b}$ gives:

$$\underline{Cv}^{n+3/2} = \underline{Cv}^{n+1/2} + \Delta t \cdot \underline{Ca}^{n+1} = \underline{b}$$

- From this we obtain:

$$\underline{Ca} = \frac{1}{\Delta t} (\underline{b} - \underline{Cv}^{n+1/2}) = \frac{1}{\bar{\gamma}} (\underline{b} - \underline{Cv}^{n+1/2}) \quad \text{having posed: } \boxed{\bar{\gamma} = \Delta t}$$

- For a variable Δt in time, one has simply:

$$\boxed{\bar{\gamma} = \frac{\Delta t^n + \Delta t^{n+1}}{2}}$$

43

Finding the Lagrange Multipliers (3)

- Summarizing, the Lagrange multipliers $\underline{\lambda}$ are obtained by solving the linear algebraic system:

$$\underline{B}^* \underline{\lambda} = \underline{w}$$

We obtain one multiplier for each imposed constraint

where the known terms are given by:

$$\underline{B}^* \equiv \underline{Cm}^{-1} \underline{C}^T \quad \text{and} \quad \underline{w} \equiv \frac{1}{\bar{\gamma}} (\underline{b} - \underline{Cv}^{n+1/2}) - \underline{Cm}^{-1} (\underline{f}^e - \underline{f}^i)$$

$$\bar{\gamma} = \Delta t \quad \text{for constant } \Delta t$$

$$\bar{\gamma} = (\Delta t^n + \Delta t^{n+1})/2 \quad \text{for variable } \Delta t \text{ in time}$$

- Finally we compute the reactions:

$$\underline{r} = \underline{C}^T \underline{\lambda}$$

We obtain one reaction for each constrained dof

and add them to the other known external forces.

- This is the only implicit part of the whole method.

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Exercise 0 – Ideal ballistics

- Motion in vacuum is analytical:

$$v_x = v_{0x} = v_0 \cos \phi_0$$

$$v_y = v_{0y} - gt = v_0 \sin \phi_0 - gt$$

- The trajectory is a parabola:

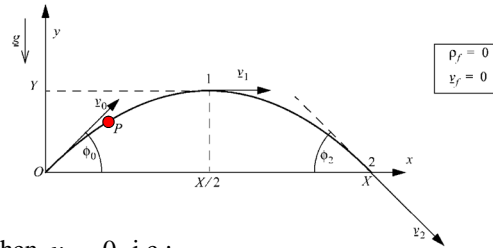
$$\phi_2 = \phi_0 \quad \text{and} \quad \|v_2\| = \|v_0\|$$

- Time to reach highest point is when $v_y = 0$, i.e.:

$$t_1 = \frac{v_{0y}}{g} = \frac{v_0}{g} \sin \phi_0$$

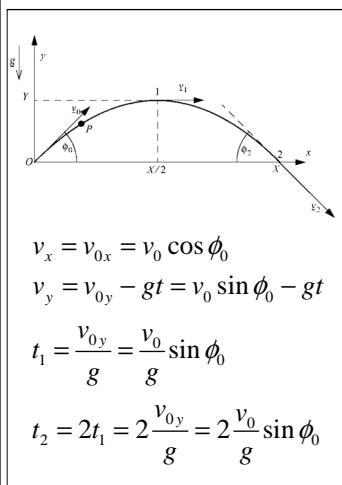
- By symmetry, time to impact is twice as long:

$$t_2 = 2t_1 = 2 \frac{v_{0y}}{g} = 2 \frac{v_0}{g} \sin \phi_0$$



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Exercise 0 – Ideal ballistics (2)



$$v_x = v_{0x} = v_0 \cos \phi_0$$

$$v_y = v_{0y} - gt = v_0 \sin \phi_0 - gt$$

$$t_1 = \frac{v_{0y}}{g} = \frac{v_0}{g} \sin \phi_0$$

$$t_2 = 2t_1 = 2 \frac{v_{0y}}{g} = 2 \frac{v_0}{g} \sin \phi_0$$

- The range is therefore:

$$X = v_x t_2 = 2v_{0x} \frac{v_{0y}}{g} = 2 \frac{v_0^2}{g} \sin \phi_0 \cos \phi_0 = \frac{v_0^2}{g} \sin(2\phi_0)$$

- Max range is when shooting at 45°:

$$X_{\max}(\phi_0) = \frac{v_0^2}{g} \quad \text{for} \quad \sin(2\phi_0) = 1 \Leftrightarrow \phi_0 = \pi/4$$

- The position at the generic time is:

$$x(t) = v_{0x} t = v_0 t \cos \phi_0$$

$$y(t) = v_{0y} t - g \frac{t^2}{2} = v_0 t \sin \phi_0 - g \frac{t^2}{2}$$

- Max elevation depends only on v_{0y} :

$$Y = y_{\max} = y(t_1) = \frac{v_{0y}^2}{g} - \frac{g}{2} \frac{v_{0y}^2}{g^2} = \frac{v_{0y}^2}{2g} = \frac{v_0^2}{2g} \sin^2 \phi_0$$

46

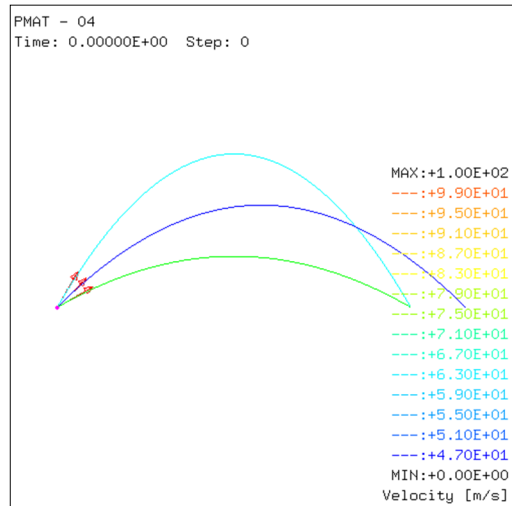
Exercise 0 – Ideal ballistics (3)

- Study motion of projectile as a function of the shooting angle:

$$v_0 = 100$$

$$\phi_0 = 30^\circ, 45^\circ, 60^\circ$$

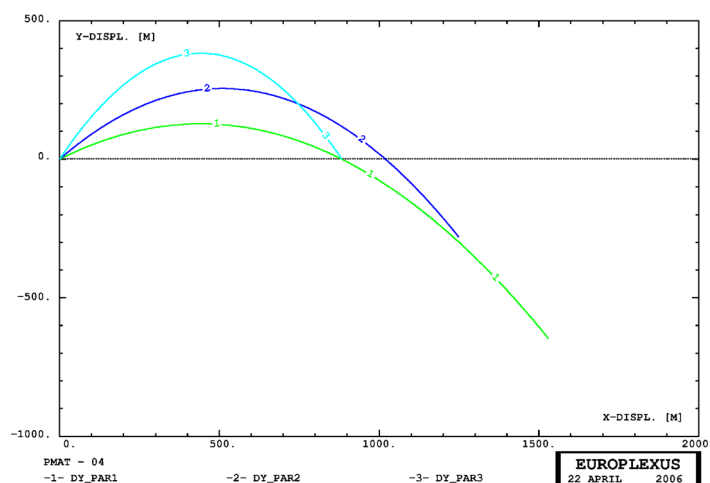
- Computed vs. analytical positions:



47

Exercise 0 – Ideal ballistics (4)

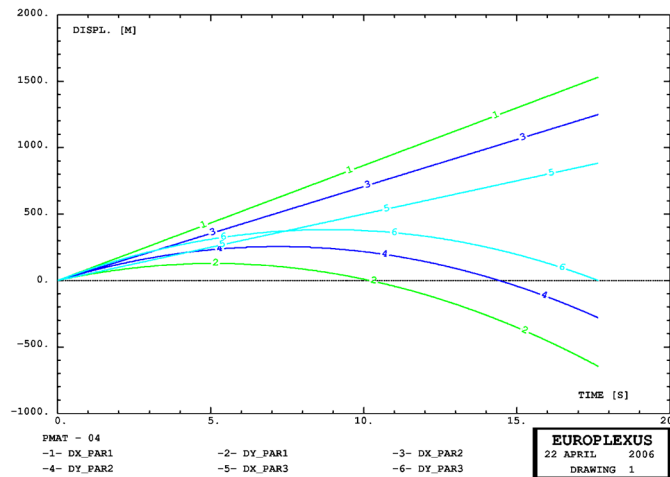
- Computed trajectories:



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Exercise 0 – Ideal ballistics (5)

- Computed displacement components:



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Exercise 0 – Ideal ballistics (6)

- Influence of time discretization:

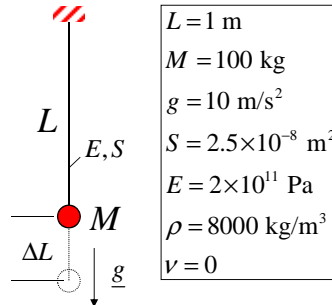
Case	Description	Flight time t_2	Range X	N. of steps
PMAT01	$\phi_0 = 30^\circ$	10.197	883.1	1000
PMAT01B	$\phi_0 = 30^\circ$	10.197	883.1	100
PMAT01C	$\phi_0 = 30^\circ$	10.197	883.1	10
PMAT01D	$\phi_0 = 30^\circ$	10.197	883.1	2
PMAT01E	$\phi_0 = 30^\circ$	10.197	883.1	1
PMAT02	$\phi_0 = 45^\circ$	14.421	1019.7	1000
PMAT03	$\phi_0 = 60^\circ$	17.662	883.1	1000
PMAT04	$\phi_0 = 30^\circ, 45^\circ, 60^\circ$	10.197, 14.421, 17.662	883.1, 1019.7	1000

- No stability restraints (no wave propagation)
- Analytical precision

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Exercise 1 – Suspended mass

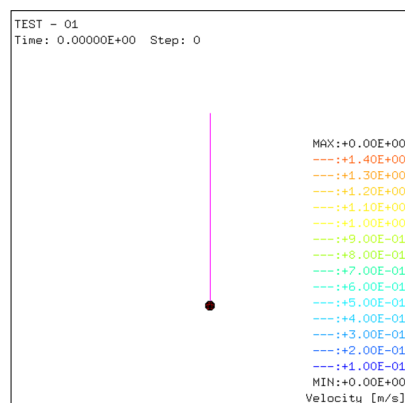
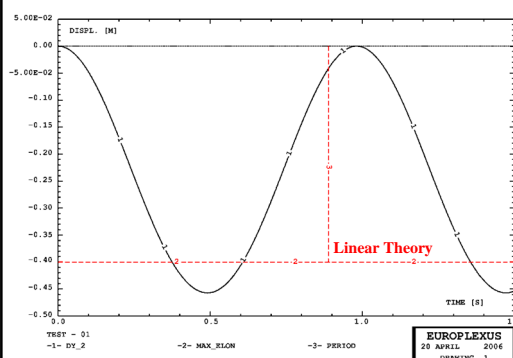
- Single-element discrete model: check vs. (linear) analytical solution
- Explain possible reasons for observed discrepancies
- Try out different values: e.g. gravity 1000 times smaller
- Discuss multi-element discrete model
- Replace gravity by initial velocity and discuss effects of structure modeling: A) as a bar, B) as a cable ...



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Exercise 1 – Suspended mass (2)

- **TEST01** : 1 element of type FUN2 (cable), no resistance to bending. Use FUNE material (no resistance to compression)



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Exercise 1 – Suspended mass (3)

- EUROPLEXUS input file:

```

TEST - 01
*-----
ECHO
* CONV win
*-----Problem type
CPLA LAOR
*-----Dimensioning
DIME
PT2L 2 FUN2 1 PMAT 1 ZONE 2
TABL 1 2
FORC 1
TERM
*-----Geometry
GEOM LIBR POIN 2 FUN2 1 PMAT 1 TERM
0 0 0 -1
1 2
2
*-----Geometrical complements
COMP EPAI 2.5E-8 LECT 1 TERM
*-----Material data
MATE FUNE RO 8000, YOUN 2.0E11 NU 0.0 ELAS 2.0E11 ERUP 1.0EO
TRAC 1 2.0E11 1.E0
LECT 1 TERM
MASS 100.0 LECT 2 TERM
*-----Boundary conditions
LINK COUP
BLOO 2 LECT 1 TERM
*-----Applied loads
CHAR 1 FACT 2 FORC 2 -1.E3 LECT 2 TERM
TABL 2 0.0 1.0 10.0 1.0
*-----Outputs
ECRI DEPL VITE CONT ECRO TFREQ 0.5
FICH ALIC TEMP FREQ 20
POIN LECT 1 2 TERM
ELEM LECT 1 TERM
*-----Options
OPTI PAS UTIL NOTEST
*-----Transient calculation
CALCUL TINI 0. TEND 1.5 PASF 0.1E-3

```

```

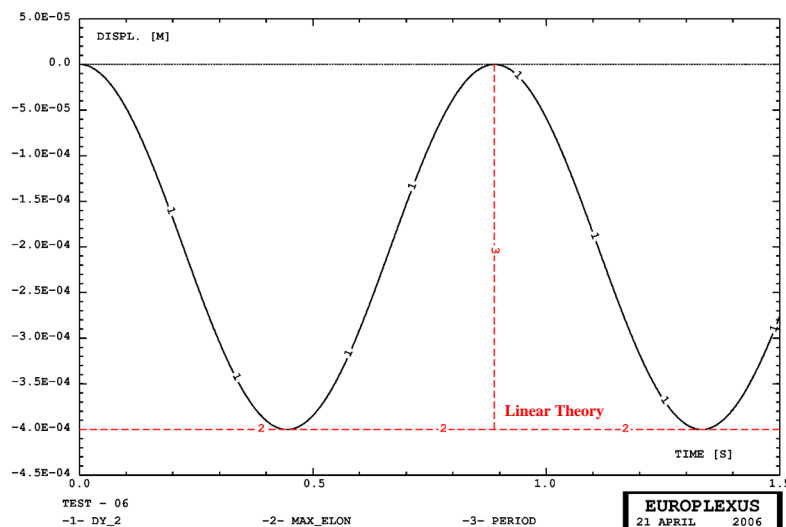
*=====POST-TREATMENT
=====
SUIT
Post-treatment
ECHO
RESU ALIC TEMP GARD PSCR
SORT GRAP
AXTE 1.0 'Time [s]'
*-----Curve definitions
COUR 1 'dy_2' DEPL COMP 2 NOEU LECT 2 TERM
COUR 2 'sg_1' CONT COMP 1 ELEM LECT 1 TERM
COUR 3 'fe_2' FORC COMP 1 NOEU LECT 2 TERM
COUR 4 'fe_1' FORC COMP 2 NOEU LECT 1 TERM
DCOU 6 'Max_elon' 2
0.0 -0.4
1.5 -0.4
DCOU 7 'Period' 2
0.889 0.0
0.889 -0.4
*-----Plots
trac 1 6 7 axes 1.0 'DISPL [M]' yzer
colo noir rouge rouge
dash 0 2 2
trac 2 axes 1.0 'CONTR. [PA]' yzer
trac 3 4 axes 1.0 'FORCE [N]' yzer
list 1 axes 1.0 'DISPL [M]'
*-----Results qualification
QUAL DEPL COMP 2 LECT 2 TERM REFE -4.53636E-1 TOLE 1.E-2
CONT COMP 1 LECT 1 TERM REFE 7.48171E+10 TOLE 1.E-2
*=====
FIN

```

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Exercise 1 – Suspended mass (4)

- TEST06 : applied load 1000 times smaller (small-strain)

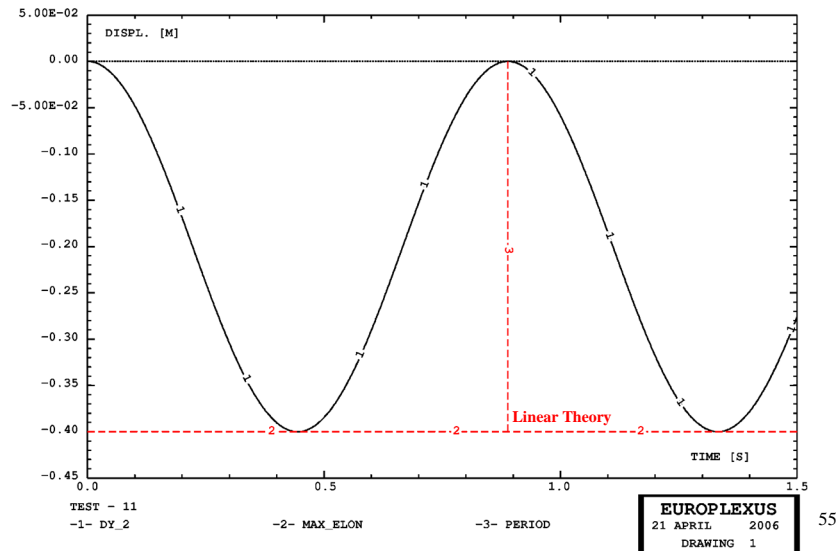


EUROPLEXUS
21 APRIL 2006
DRAWING 1

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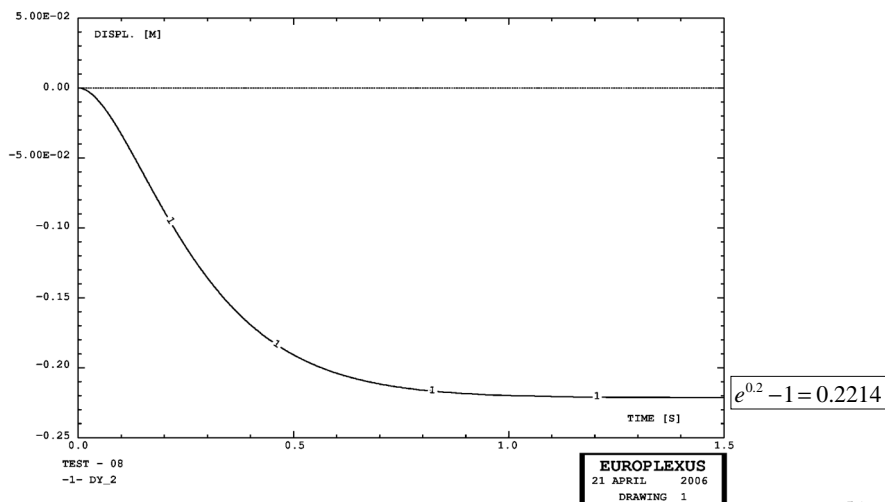
Exercise 1 – Suspended mass (5)

- **TEST11** : standard load but small-strain option (OPTI EDSS)



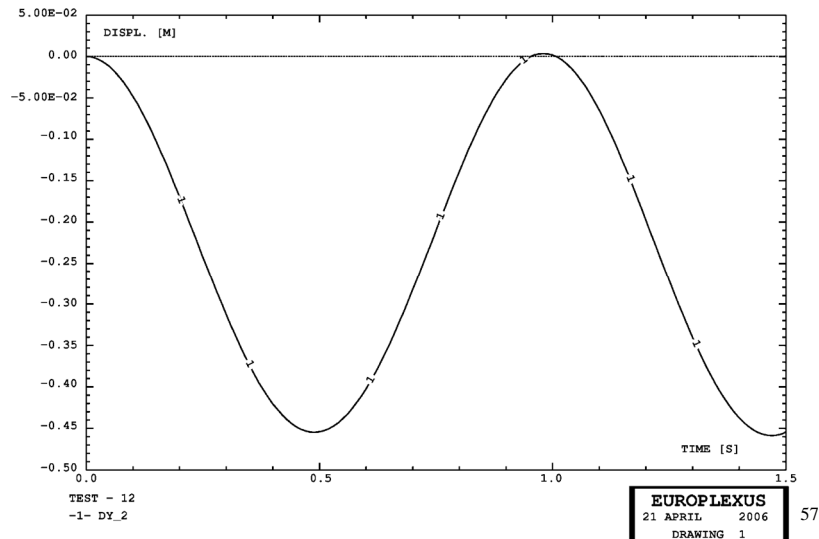
Exercise 1 – Suspended mass (6)

- **TEST08** : “quasi-static” option (OPTI QUAS STAT ...)



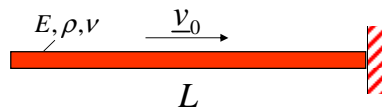
Exercise 1 – Suspended mass (7)

- **TEST12** : time increment 50 times the critical value



Exercise 2 – Wave propagation

- Obtain 1D analytical solution
- Discuss numerical solution
- Why was cross-section not specified?
- Study effect of time increment
- Study effect of Poisson's ratio ...

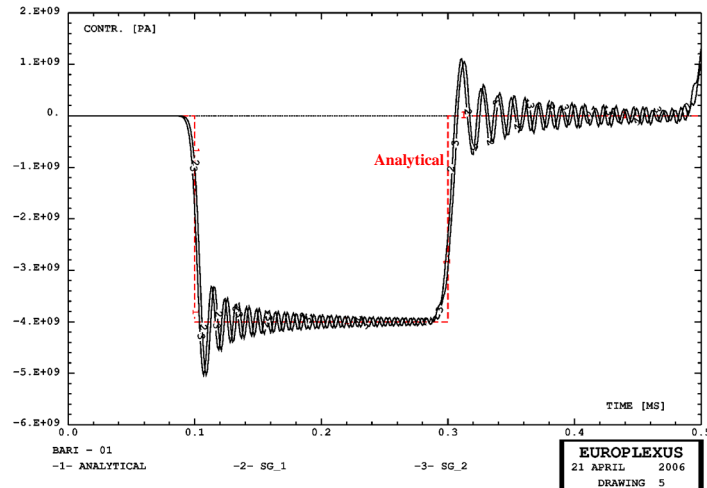


$L = 1 \text{ m}$
$v_0 = 100 \text{ m/s}$
$E = 2 \times 10^{11} \text{ Pa}$
$\rho = 8000 \text{ kg/m}^3$
$\nu = 0.3$

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Exercise 2 – Wave propagation (2)

- **BARI01** : 100 elements of type FUN2 (cable), VM23 material (traction/compression), small-strain option (OPTI EDSS)



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Exercise 2 – Wave propagation (3)

- EUROPLEXUS input file:

```

BARI - 01
*-----
ECHO
*CONV win
CAST mesh
*-----Problem type
CPLA LAGR LAGC
*-----Dimensioning
DIME
PT2L 102 FUN2 100 PMAT 1 ZONE 2
BLOQ 2
TABL 1 2
FORC 1
IMPA 1 PSIM 1
TERM
*-----Geometry
GEOM FUN2 bar PMAT obst TERM
*-----Geometrical complements
COMP EPAI 2.5E-8 LECT bar TERM
*-----Material data
MATE VM23 RO 8000. YOUN 2.0E11 NU 0.0 ELAS 2.0E11
TRAC 1 2.0E11 1.E0
LECT bar TERM
MASS 1.0 LECT obst TERM
*-----Boundary conditions
LIAI FREQ 1
BLOQ 12 LECT obst TERM
IMPA DDL 1 COTE -1
PROJ LECT obst TERM
CIBL LECT p2 TERM
*-----Initial conditions
INIT VITE 1 100.0 LECT bar TERM
*-----Outputs
ECRI DEPL VITE CONT ECRO TFREQ 0.1E-3
FICH ALIC TEMP FREQ 1
POIN LECT p3 p2 TERM
ELEM LECT e1 e2 TERM
*-----Options
OPTI PAS UTIL NOTEST
LOG 1
EDSS
*-----Transient calculation
CALCUL TINI 0. TEND 0.5E-3 PASF 0.1E-5
    
```

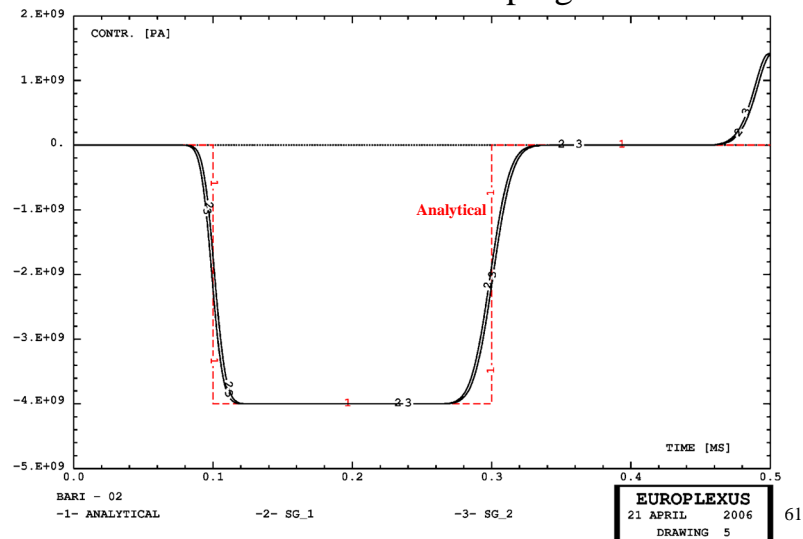
```

*=====POST-TREATMENT
=====
SUIT
Post-treatment
ECHO
RESU ALIC TEMP GARD PSCR
SORT GRAP
AXTE 1000.0 'Time [ms]'
*-----Curve definitions
COUR 1 'dx_2' DEPL COMP 1 NOEU LECT p2 TERM
COUR 2 'dx_3' DEPL COMP 1 NOEU LECT p3 TERM
COUR 3 'sg_1' CONT COMP 1 ELEM LECT e1 TERM
COUR 4 'sg_2' CONT COMP 1 ELEM LECT e2 TERM
DCOU 5 'Analytical' 6
0.0 0.0
0.1E-3 0.0
0.1E-3 -4.E9
0.3E-3 -4.E9
0.3E-3 0.0
0.5E-3 0.0
*-----Plots
trac 1 axes 1.0 'DISPL. [M]'
trac 2 axes 1.0 'DISPL. [M]'
trac 3 axes 1.0 'CONTR. [PA]'
trac 4 axes 1.0 'CONTR. [PA]'
trac 5 3 4 axes 1.0 'CONTR. [PA]' yzer
colo roug noir noir
dash 2 0 0
*-----Results qualification
QUAL DEPL COMP 1 LECT p2 TERM REFE -9.76759E-3 TOLE 1.E-2
DEPL COMP 1 LECT p3 TERM REFE -9.86359E-3 TOLE 1.E-2
CONT COMP 1 LECT e1 TERM REFE 1.00128E+9 TOLE 1.E-2
CONT COMP 1 LECT e2 TERM REFE 1.31288E+9 TOLE 1.E-2
*=====
FIN
    
```

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Exercise 2 – Wave propagation (4)

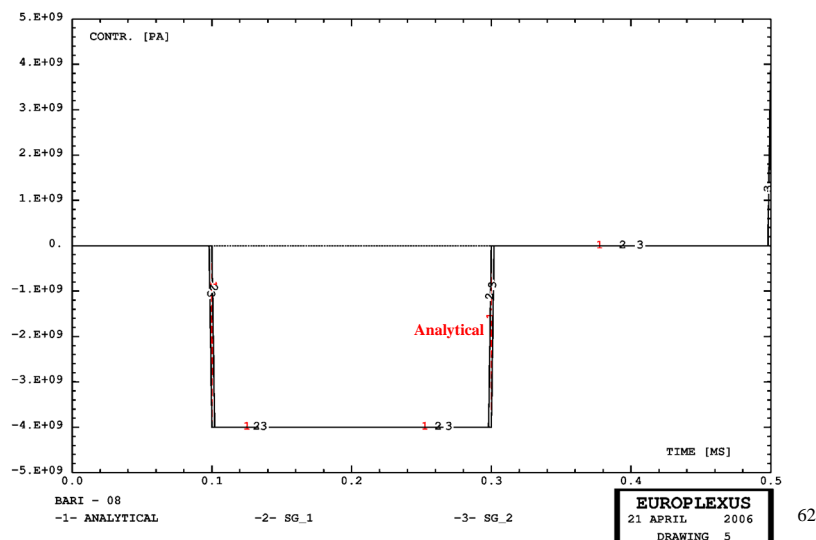
- **BARI02** : 10 % of critical damping added



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Exercise 2 – Wave propagation (5)

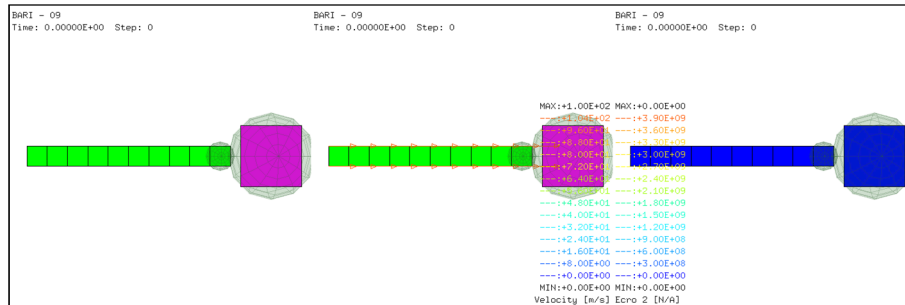
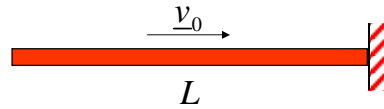
- **BARI08** : use critical time increment



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Exercise 2 – Wave propagation (6)

- **BARI09** : 2-D geometry (Q42L) and pinballs for contacts:



Geometry

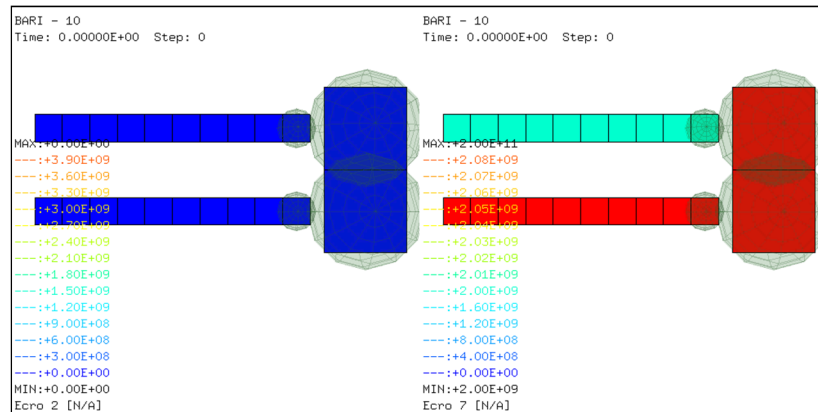
Velocities

Von Mises

63

Exercise 2 – Wave propagation (7)

- **BARI10** : Compare elastic (bottom) and elasto-plastic (top) solutions:



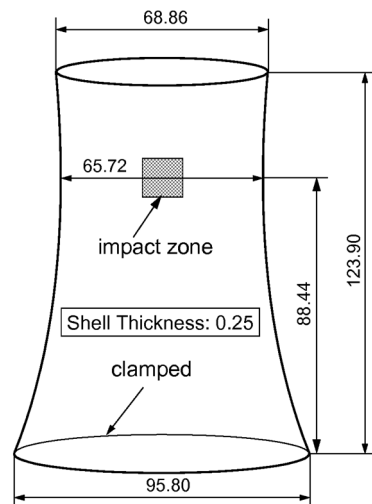
Von Mises

Yield Limit

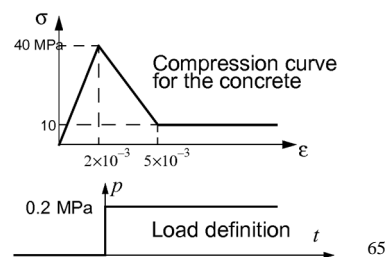
64

Exercise 3 – Impact on Cooling Tower

- Problem definition:

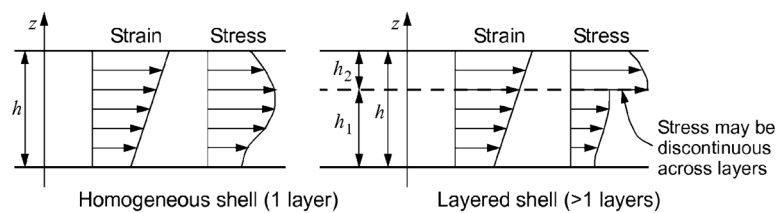


Layer N.	ϕ_L	Material
1	0.1	concrete
2	0.0025	steel
3	0.795	concrete
4	0.0025	steel
5	0.1	concrete

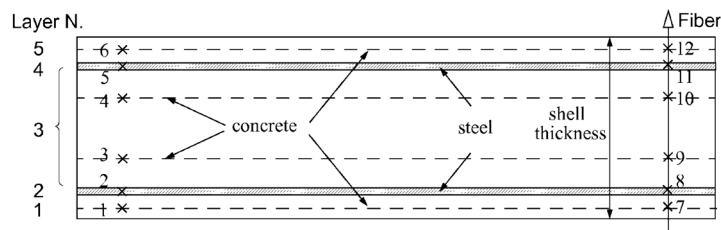


Exercise 3 – Impact on Cooling Tower (2)

- Shell model:

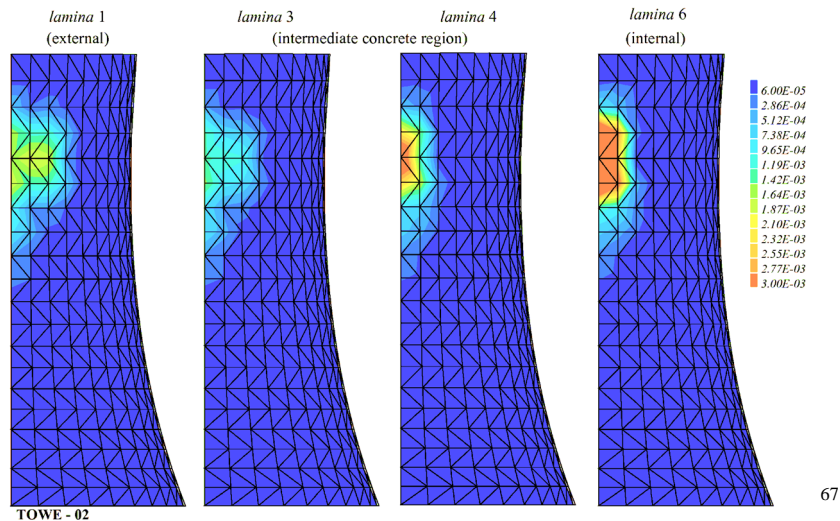


- Layers:

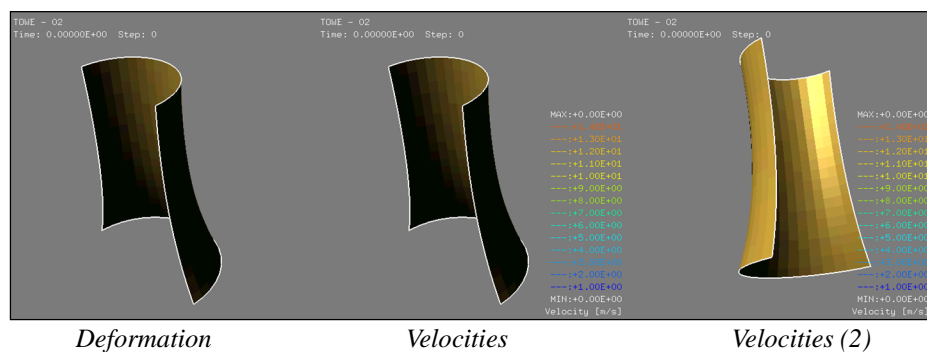


Exercise 3 – Impact on Cooling Tower (3)

- Results - equivalent plastic strain for the 4 concrete laminae at 100 ms:



Exercise 3 – Impact on Cooling Tower (4)

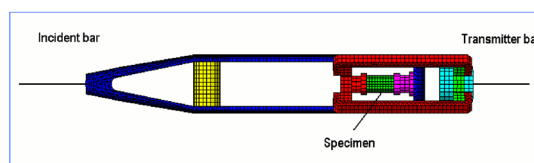


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Exercise 4 – Hopkinson Bar



JRC Large Dynamic Test Facility (HopLab)

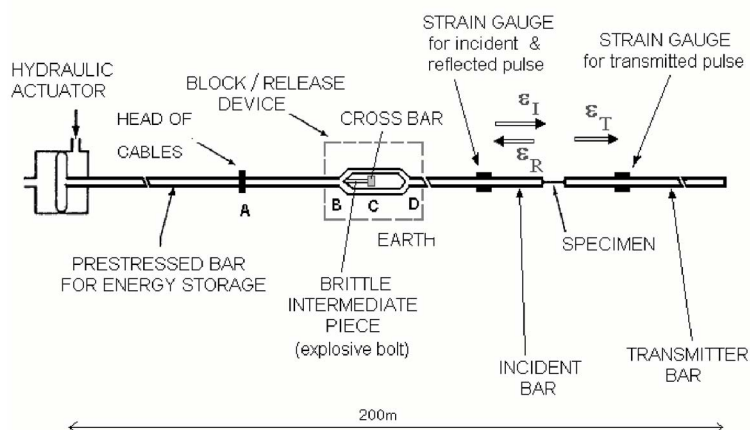


Fast Compression Test on Ductile Metallic Specimen

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Exercise 4 – Hopkinson Bar (2)

Functioning Principle (Hopkinson):



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Exercise 4 – Hopkinson Bar (3)

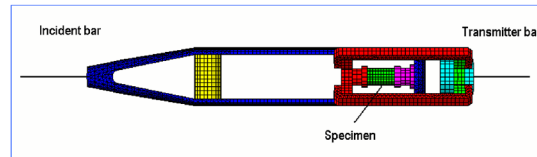
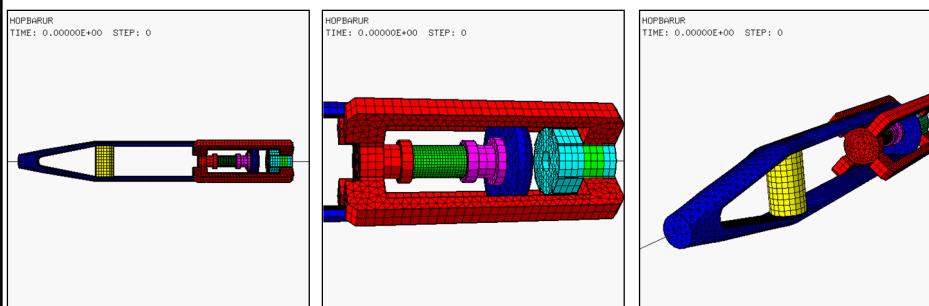


Figure 2: Incident and transmitter inversion cages

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Exercise 4 – Hopkinson Bar (4)

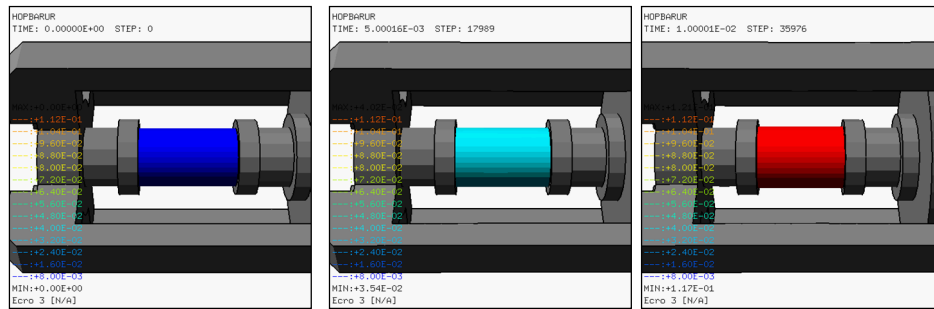
Numerical simulation : geometrical model



72

Exercise 4 – Hopkinson Bar (5)

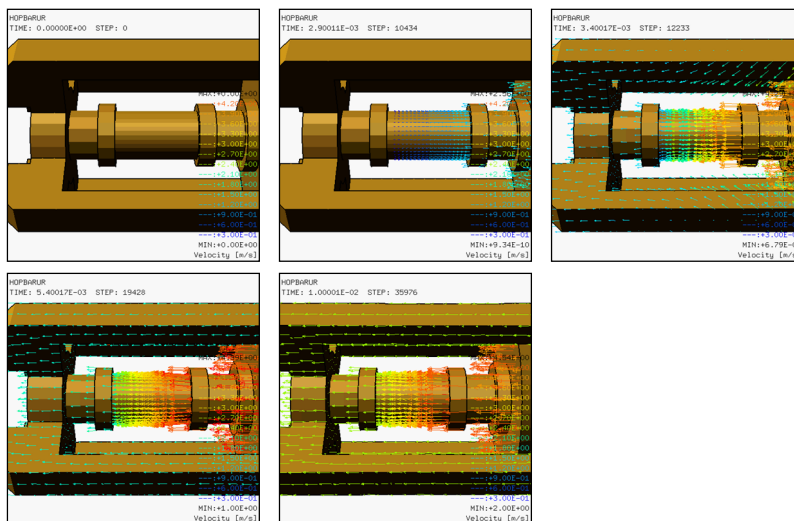
Numerical simulation : specimen plastification



73

Exercise 4 – Hopkinson Bar (6)

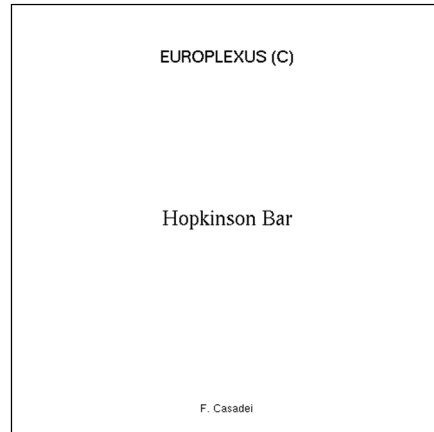
Numerical simulation : velocities



74

Exercise 4 – Hopkinson Bar (7)

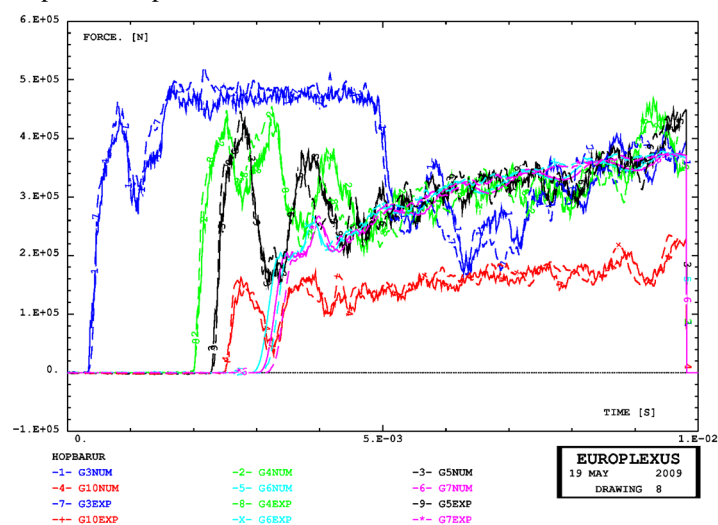
Numerical simulation : results animation



75

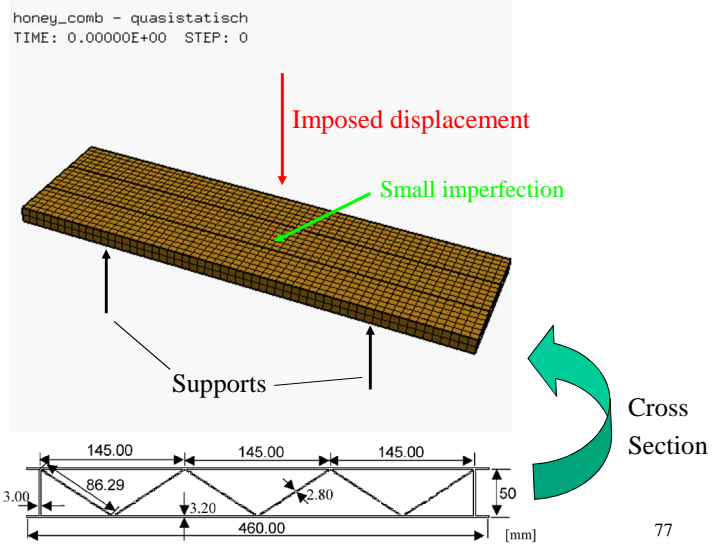
Exercise 4 – Hopkinson Bar (8)

Comparison experiment vs. calculation :

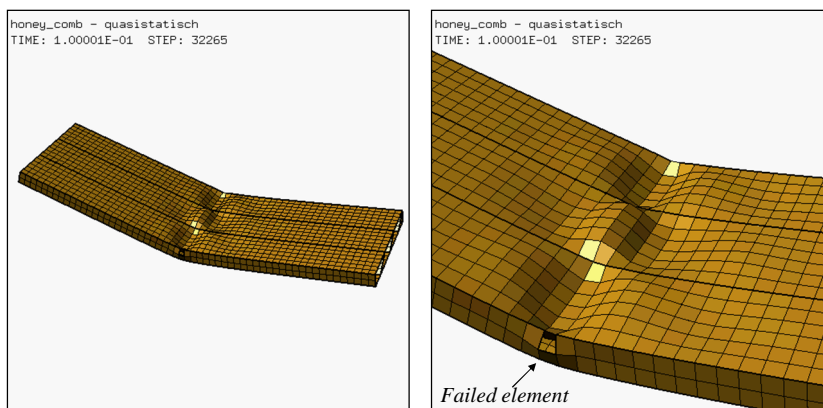


76

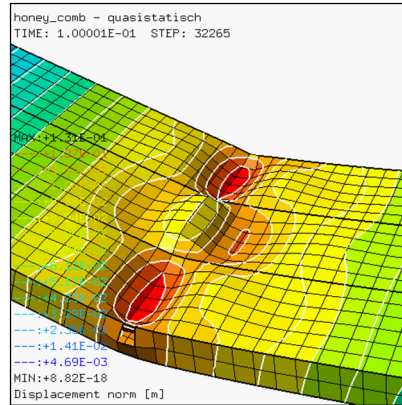
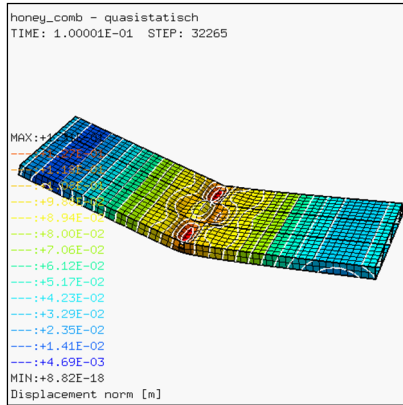
Exercise 5 – Honeycomb buckling



Exercise 5 – Honeycomb buckling (2)



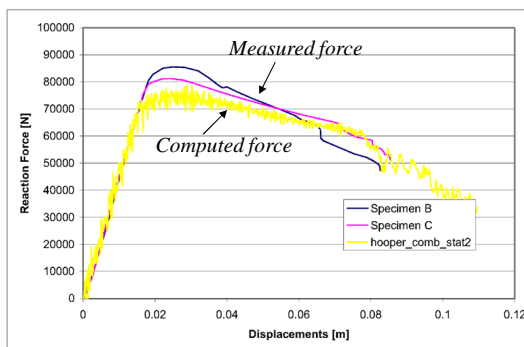
Exercise 5 – Honeycomb buckling (3)



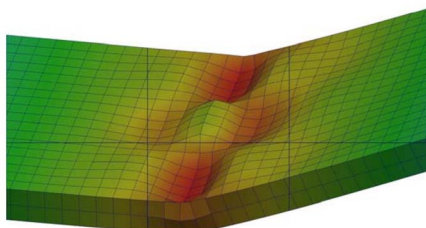
Final displacement norm (buckling mode)

79

Exercise 5 – Honeycomb buckling (4)

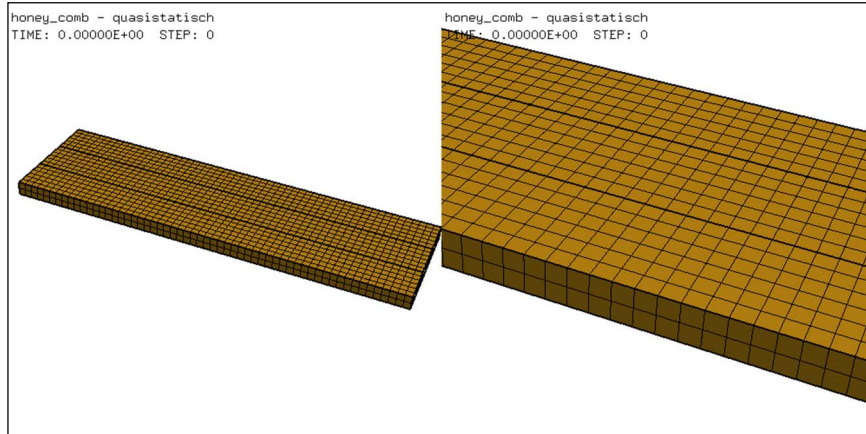


Comparison
experiment /
calculation



80

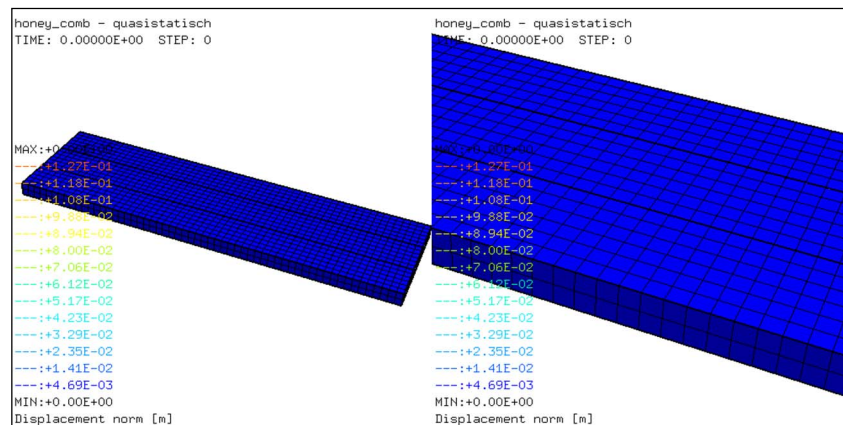
Exercise 5 – Honeycomb buckling (5)



Deformation

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Exercise 5 – Honeycomb buckling (6)



Displacement

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